# CAE 208 / MMAE 320: Thermodynamics Fall 2023

# **November 21, 2023** Entropy (3)

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# ANNOUNCEMENTS

- Exam 2 solutions are posted
- Assignment 9 is posted (2-3 days more extension is fine to submit it on Friday or the weekend after the Thanksgiving)
- I will post assignment 10 due 12/01/23 (those who need to submit it)
- The bonus activities document is posted. Please pay attention to the deadlines and also the updates about this task (an idea submission by the 11/29/2023 is required)

## RECAP

 The equality in the Clausius inequality holds for totally or jut internally reversible cycles and the inequality for the irreversible ones:

$$\left(\oint \frac{\delta Q}{T}\right)_{int,rev} = 0$$

#### Recap

A special case: Internally reversible isothermal heat transfer processes:

$$\Delta S = \int_{1}^{2} \left(\frac{\delta Q}{T}\right)_{int,rev} = \int_{1}^{2} \left(\frac{\delta Q}{T_{0}}\right)_{int,rev} = \frac{1}{T_{0}} \int_{1}^{2} \delta Q_{int,rev}$$

$$\Delta S_{isothermal} = \frac{Q}{T_0} \qquad \left(\frac{kJ}{K}\right)$$

$$T = 300 \text{ K} = \text{const.}$$
$$\Delta S_{\text{sys}} = \frac{Q}{T} = 2.5 \frac{\text{kJ}}{\text{K}}$$
$$Q = 750 \text{ kJ}$$

• Entropy is a property:



• For entropy, we can say "the increase of entropy principle":

$$\oint_1^2 \frac{\delta Q}{T} \leq \Delta S$$

$$\Delta S_{sys} = S_2 - S_1 = \oint_1^2 \frac{\delta Q}{T} + S_{gen}$$

 $\Delta S_{isolated} \geq 0$ 

• For an isolated system:

$$S_{gen} = \Delta S_{total} = \Delta S_{sys} + \Delta S_{surr} \ge 0$$

 $S_{gen} = \begin{cases} > 0. & irreversible process \\ = 0 & reversible process \\ < 0. & impossible process \end{cases}$ 



#### Recap

- The entropy of a fixed mass can be changed by:
   Heat Transfer
   Irreversibilities
- Entropy of a fixed mass does not change during a process that is internally reversible and adiabatic. During this process entropy remains constant and we call it *isentropic* process

$$\Delta s = 0 \quad or \quad s_2 = s_1 \quad \left(\frac{kJ}{kg - K}\right)$$



• We can rearrange our entropy equation:

$$\delta Q_{int,rev} = T \, dS \qquad \qquad Q_{int,rev} = \int_{1}^{2} T \, dS \qquad (kJ)$$



(The area under the process curve on a T-S diagram represents heat transfer during an internally reversible process)

# WHAT IS ENTROPY

 Entropy can be viewed as a measure of molecular disorder or molecular randomness (lowest in solid phase and highest in the gas phase)



- Oscillations in solid phase fade as temperature is decreased, and the molecules supposedly become motionless at absolute zero. This represents a state of ultimate molecular order (and *minimum energy*).
- The entropy of a pure crystalline substance at absolute zero temperature is zero since there is no uncertainty about the state of the molecules at that instant (3rd law of thermodynamics)



- The third law of thermodynamics provide an absolute reference point for the determination of entropy
- The entropy determined relative to this point is called absolute entropy and it is extremely useful in the thermodynamics analysis of entropy
- The entropy of a substance that is not pure crystalline (such as a solid solution) is not zero at absolute zero temperature

• A few examples:





# THE T DS RELATIONS

#### The T ds Relations

The first T ds (or Gibbs) equation (closed stationary system):

$$\delta Q_{int,rev} - \delta W_{int,rev,out} = dU$$

 $TdS = dU + PdV \qquad (kJ)$ 

 $\delta Q_{int,rev} = Tds$ 

$$Tds = du + Pdv \quad (\frac{kJ}{kg})$$

 $\delta W_{int,rev,out} = PdV$ 

$$Tds = du + Pdv$$
  

$$h = u + Pv \rightarrow dh = du + Pdv + vdP \qquad \Rightarrow Tds = dh - vdP$$

#### The T ds Relations

Use the first T ds (or Gibbs) equation to solve for entropy changes

$$ds = \frac{du}{T} + \frac{Pdv}{T}$$

$$ds = \frac{dh}{T} - \frac{vdP}{T}$$

Closed  
system
$$CV$$
 $T ds = du + P dv$   
 $T ds = dh - v dP$ 

# ENTROPY CHANGE OF LIQUIDS AND SOLIDS

#### **Entropy Change of Liquids and Solids**

 Liquids and solids can be approximated as incompressible substances (*dv* ≅ 0 & *c*<sub>p</sub> = *c*<sub>v</sub> = *c*<sub>p</sub> = *c*):

$$ds = \frac{dh}{T} - \frac{vdP}{T}$$

$$s_2 - s_1 = \int_1^2 c(T) \frac{dT}{T} \cong c_{avg} \ln(\frac{T_2}{T_1})$$

$$s_2 - s_1 = \int_1^2 c(T) \frac{dT}{T} \cong c_{avg} \ln\left(\frac{T_2}{T_1}\right) = 0 \quad \rightarrow \quad T_2 = T_1$$
 (For isentropic)

# THE ENTROPY CHANGE OF IDEAL GASES

• For gases, we can write:

$$ds = \frac{du}{T} + \frac{Pdv}{T} = c_v \frac{dT}{T} + R \frac{dv}{v}$$

$$s_2 - s_1 = \int_1^2 c_v(T) \frac{dT}{T} + R \ln(\frac{v_2}{v_1})$$

$$s_2 - s_1 = \int_1^2 c_p(T) \frac{dT}{T} + R \ln(\frac{P_2}{P_1})$$

• Approach 1: Constant Specific Heats (Approximate Analysis):

$$s_{2} - s_{1} = \int_{1}^{2} c_{v}(T) \frac{dT}{T} + R \ln(\frac{v_{2}}{v_{1}})$$

$$s_{2} - s_{1} = c_{v,avg} \ln\left(\frac{T_{2}}{T_{1}}\right) + R \times ln\left(\frac{v_{2}}{v_{1}}\right)$$

$$c_{p,avg}$$

$$c_{p,avg}$$

$$C_{p,avg}$$

$$T_{1}$$

$$T_{avg}$$

$$T_{2}$$

$$T_{1}$$

$$\begin{cases} dh = c_p dT \\ v = \frac{RT}{P} \end{cases} \rightarrow s_2 - s_1 = c_{p,avg} \times \ln\left(\frac{T_2}{T_1}\right) - R \times \ln\left(\frac{P_2}{P_1}\right) \end{cases}$$

• Approach 2: Variable Specific Heats (Exact Analysis):

$$s^{0} = \int_{0}^{T} c_{p}(T) \frac{dT}{T}$$
$$\int_{0}^{T} c_{p}(T) \frac{dT}{T} = s_{2}^{0} - s_{1}^{0}$$
$$s_{2} - s_{1} = s_{2}^{0} - s_{1}^{0} - R \times \ln(\frac{P_{2}}{P_{1}})$$
$$\bar{s}_{2} - \bar{s}_{1} = \bar{s}_{2}^{0} - \bar{s}_{1}^{0} - R_{u} \times \ln(\frac{P_{2}}{P_{1}})$$

• Approach 2: Variable Specific Heats (Exact Analysis):



# **CLASS ACTIVITY**

- Air is compressed from an initial state of 100 kPa and 17 °C to a final state of 600 kPa and 57 °C. Determine the entropy change of air during this compression process by using:
  - a) Property values from the air table
  - b) Average specific heats

 Solution (assumptions):
 Air is an ideal gas since it is at a high temperature and low pressure relative to its critical point values



Solution (calculations):
 Part (a): Table A-21

$$s_{2} - s_{1} = s_{2}^{0} - s_{1}^{0} - R \times \ln\left(\frac{P_{2}}{P_{1}}\right) = \left[(1.79783 - 1.66802)\frac{kJ}{kg - K}\right] \times \ln(\frac{600 \, kPa}{100 \, kPa})$$
$$s_{2} - s_{1} = -0.3844 \frac{kJ}{kg - K}$$

Part (b): Using a c<sub>p</sub> value at the average temperature of 37 °C (Table A-2b)

$$s_{2} - s_{1} = c_{p,avg} \ln\left(\frac{T_{2}}{T_{1}}\right) - R \times \ln\left(\frac{P_{2}}{P_{1}}\right)$$
$$= \left(1.006 \frac{kJ}{kg - K}\right) \ln\left(\frac{330 K}{290 K}\right) - \left(0.287 \frac{kJ}{kg - K}\right) \ln\left(\frac{600 kPa}{100 kPa}\right) = -0.3842 \frac{kJ}{kg - K}$$

# ISENTROPIC PROCESSES OF IDEAL GASES

#### **Isentropic Processes of Ideal Gases**

 Approach 1: Constant Specific Heats (Approximate Analysis) for Isentropic Processes of Ideal Gases

$$s_2 - s_1 = 0 = c_v \ln\left(\frac{T_2}{T_1}\right) + R \times ln\left(\frac{v_2}{v_1}\right)$$

$$\ln\left(\frac{T_2}{T_1}\right) = -\frac{R}{c_v}\ln\left(\frac{v_2}{v_1}\right) \qquad \begin{cases} c_p - c_v = R\\ k = \frac{c_p}{c_v} & \rightarrow \frac{R}{c_v} = k - 1 \end{cases}$$

$$\ln\left(\frac{T_2}{T_1}\right) = \ln\left(\frac{v_1}{v_2}\right)^{\frac{R}{c_v}}$$

$$\frac{T_2}{T_1} = \left(\frac{v_1}{v_2}\right)^{k-1}$$

#### **Isentropic Processes of Ideal Gases**

 Approach 1: Constant Specific Heats (Approximate Analysis) for Isentropic Processes of Ideal Gases

$$\left(\frac{T_2}{T_1}\right)_{s=constant} = \left(\frac{v_1}{v_2}\right)^{k-1}$$

$$\left(\frac{T_2}{T_1}\right)_{s=constant} = \left(\frac{P_2}{P_1}\right)^{\frac{k-1}{k}}$$

$$\left(\frac{P_2}{P_1}\right)_{s=constant} = \left(\frac{v_1}{v_2}\right)^k$$

#### **Isentropic Processes of Ideal Gases**

 Approach 1: Constant Specific Heats (Approximate Analysis) for Isentropic Processes of Ideal Gases

 $Tv^{k-1} = Constant$   $\left(\frac{T_2}{T_1}\right)_{s=const.} = \left(\frac{P_2}{P_1}\right)^{(k-1)/k} = \left(\frac{v_1}{v_2}\right)^{k-1}$   $TP^{\frac{1-k}{k}} = Constant$   $Valid for \quad *ideal gas \\*isentropic process \\*constant specific heats}$ 

 Approach 2: Variable Specific Heats (Exact Analysis) for Isentropic Processes of Ideal Gases

$$0 = s_2^0 - s_1^0 - R \times \ln(\frac{P_2}{P_1})$$

$$s_2^0 = s_1^0 + R \times \ln(\frac{P_2}{P_1})$$

$$s_2^0 = s_1^0 + R \times \ln\left(\frac{P_2}{P_1}\right) \rightarrow \frac{P_2}{P_1} = \exp(\frac{s_2^0 - s_1^0}{R})$$

# **CLASS ACTIVITY**

 Air enters an isentropic turbine at 150 psia and 900 °F through a 0.5 ft<sup>2</sup> inlet section with a velocity of 500 ft/s. It leaves at 15 psia with a velocity of 100 ft/s. Calculate the air temperature at the turbine exit and the power produced, in hp, by this turbine.

- Solution (assumptions):
  - □ Steady flow
  - □ The process is isentropic (both reversible and adiabatic)
  - □ Ideal gas with a constant specific heat



• Solution (Tables):  $\Box$  Table A-2Eb: @600 °F  $\rightarrow c_p = 0.250 \frac{Btu}{lbm-R}$  and k = 1.3777

Temp., °F	$c_p$ Btu/lbm $\cdot  \mathbb{R}$	$c_{\rm v}$ Btu/lbm $\cdot{\rm R}$	k	$c_p$ Btu/lbm $\cdot {\rm R}$	$c_{v}$ Btu/lbm $\cdot$ R	k
		Air	Carbon dioxide, CO <sub>2</sub>			
40	0.240	0.171	1.401	0.195	0.150	1.300
100	0.240	0.172	1.400	0.205	0.160	1.283
200	0.241	0.173	1.397	0.217	0.172	1.262
300	0.243	0.174	1.394	0.229	0.184	1.246
400	0.245	0.176	1.389	0.239	0.193	1.233
500	0.248	0.179	1.383	0.247	0.202	1.223
600	0.250	0.182	1.377	0.255	0.210	1.215
700	0.254	0.185	1.371	0.262	0.217	1.208
800	0.257	0.188	1.365	0.269	0.224	1.202
900	0.259	0.191	1.358	0.275	0.230	1.197
1000	0.263	0.195	1.353	0.280	0.235	1.192
1500	0.276	0.208	1.330	0.298	0.253	1.178
2000	0.286	0.217	1.312	0.312	0.267	1.169

• Solution (Tables):  $\Box$  Table A-1E:  $R = 0.3704 \frac{psia-ft^3}{lbm-R}$ 

TABLE A-1E											
Molar mass, gas constant, and critical-point properties											
		Molar mass, <i>M</i> lbm/lbmol	Gas constant, R*		Critical-point properties						
Substance	Formula		Btu/lbm · R	psia · ft <sup>3</sup> /lbm · R	Temperature, R	Pressure, psia	Volume, ft <sup>3</sup> /lbmol				
Air	_	28.97	0.06855	0.3704	238.5	547	1.41				

• Solution (Problem solving):

 $\dot{m} = \dot{m}_1 = \dot{m}_2$ 

$$\dot{E}_{in} - \dot{E}_{out} = \frac{d\dot{E}_{system}}{dt} = 0$$

$$\dot{m}(h_1 + V_1^2) = \dot{m}\left(h_2 + \frac{V_2^2}{2}\right) + \dot{W}_{out}$$

$$\dot{W}_{out} = \dot{m} \left( h_1 - h_1 + \frac{V_1^2 - V_2}{2} \right)$$



• Solution (Calculations):

$$\frac{T_2}{T_1} = \left(\frac{P_2}{P_1}\right)^{\frac{k-1}{k}} \to T_2 = T_1 \times \left(\frac{P_2}{P_1}\right)^{\frac{k-1}{k}} = (900 + 460 R) \left(\frac{15 \text{ psia}}{150 \text{ psia}}\right)^{\frac{0.3777}{1.377}} = 724 R$$

$$v_1 = \frac{RT_1}{P_1} = \frac{\left(0.3704\frac{psia - ft^3}{lbm - R}\right)(900 + 460\,R)}{150\,psia} = 3.358\frac{ft^3}{lbm}$$

$$\dot{m} = \frac{A_1 V_1}{v_1} = \frac{(0.5 ft^2) \left(500 \frac{ft}{s}\right)}{3.358 \frac{ft^3}{lbm}} = 74.45 \frac{lbm}{s}$$

• Solution (Calculations):

$$\dot{W}_{out} = \dot{m} \left( h_1 - h_1 + \frac{V_1^2 - V_2}{2} \right)$$

$$\dot{W}_{out} = \left(74.45 \frac{lbm}{s}\right) \left[ \left(0.250 \frac{Btu}{lbm - R}\right) (1360 - 724R) + \left(\frac{\left(500 \frac{ft}{s}\right)^2}{2} - \frac{\left(100 \frac{ft}{s}\right)^2}{2}\right) \left(\frac{1 \frac{Btu}{lbm}}{25.037} ft^2 - \frac{100 \frac{ft}{s}}{s^2}\right) \right] \left(\frac{1 \frac{Btu}{lbm}}{s^2} - \frac{100 \frac{ft}{s}}{s^2}\right) \left(\frac{1 \frac{Btu}{lbm}}{s^2} - \frac{100 \frac{ft}$$

$$\dot{W}_{out} = 12,194 \frac{Btu}{s} \left( \frac{1 \ hp}{0.7068 \frac{Btu}{s}} \right) = 17,250 \ hp$$

## **REVERSIBLE STEADY-FLOW WORK**

Recall we had

$$W_b = \int_1^2 P \ dV$$

• For steady flow:

$$\delta q_{rev} - \delta w_{rev} = dh + dke + dpe$$

$$-\delta w_{rev} = vdP + dke + dpe$$

 $\begin{cases} \delta q_{rev} = Td \\ Tds = dh - vdP \end{cases} \rightarrow \delta q_{rev} = dh - vdP$ 

$$w_{rev} = -\int_{1}^{2} v dP - \Delta ke - \Delta pe$$

#### **Reversible Steady-Flow Work**

• For steady flow:

$$w_{rev} = -\int_{1}^{2} v dP - \Delta ke - \Delta pe$$



(a) Steady-flow system



#### **Reversible Steady-Flow Work**

• We write the *Bernoulli equation*:

$$w_{rev} = 0 = v(P_2 - P_1) + \left(\frac{V_2^2 - V_1^2}{2}\right) + g(z_2 - z_1)$$

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# **ISENTROPIC EFFICIENCIES OF STEADY-FLOW DEVICES**

 Steady-flow devices deliver the most and consume the least work when the process is reversible:



#### **Isentropic Efficiencies**

• Isentropic efficiency of a turbine can be written as:

$$\eta_T = \frac{Actual \ turbine \ work}{Isentropic \ turbine \ work} = \frac{w_a}{w_s}$$

$$\eta_T \cong \frac{h_1 - h_{2a}}{h_1 - h_{2s}}$$





Isentropic efficiency of compressors and pumps

$$\eta_{C} = \frac{Isentropic\ compressor\ work}{Actual\ compressor\ work} = \frac{w_{s}}{w_{a}}$$

$$\eta_C \cong \frac{h_{2s} - h_1}{h_{2a} - h_1}$$



#### **Isentropic Efficiencies**

Isentropic efficiency of nozzles

$$\eta_N = \frac{Actual \ KE \ nozzle \ exit}{Isentropic \ KE \ at \ nozzle \ exit} = \frac{V_{2a}^2}{V_{2s}^2}$$

$$h_1 = h_{2a} + \frac{V_{2a}^2}{2}$$

$$\eta_N \cong \frac{h_1 - h_{2a}}{h_1 - h_{2s}}$$



# **CLASS ACTIVITY**

- Steam enters an adiabatic turbine steadily at 3 MPa and 400 C and leaves at 50 kPa and 100 C. If the power output of the turbine is 2 MW, determine:
  - a) The isentropic efficiency of the turbine
  - b) The mass flow rate of the steam flowing through the turbine

- Solution (assumptions):
  - □ Steady operating conditions exist
  - □ The kinetic and potential energies are negligible





• Solution (Tables):

State 1: 
$$\begin{cases} P_1 = 3 MPa \\ T_1 = 400 \ ^{\circ}C \end{cases} \rightarrow \begin{cases} h_1 = 3231.7 \ kJ/kg \\ s_1 = 6.9235 \ kJ/(kg - K) \end{cases}$$

State 2a: 
$$\begin{cases} P_1 = 50 \ kPa \\ T_{2a} = 100 \ ^{\circ}C \end{cases} \rightarrow h_{2a} = 2682.4 \ kJ/kg$$

State 2s: 
$$\begin{cases} P_1 = 50 \ kPa \\ s_1 = s_2 \end{cases} \rightarrow \begin{cases} s_f = 1.0912 \ kJ/(kg - K) \\ s_g = 7.5931 \ kJ/(kg - K) \end{cases}$$

$$x_{2s} = \frac{s_{2s} - s_f}{s_{fg}} = \frac{6.9235 - 1.0912}{6.5019} = 0.897$$

$$h_{2s} = h_f + x_{2s} \times h_{fg} = 340.54 + 0.897 \times (2304.7) = 2407.9 \, kJ/kg$$

• Solution (a): The isentropic efficiency is:

$$\eta_T = \frac{h_1 - h_{2a}}{h_1 - h_{2s}} = \frac{3231.7 - 2682.4}{3231.7 - 2407.9} = 0.667 \ (or \ 66.7\%)$$

• Solution (b): The mass flow rate is:

$$\begin{split} \dot{E}_{in} &= \dot{E}_{out} \\ \dot{m}h_1 &= \dot{W}_{a,out} + \dot{m}h_{2a} \\ \dot{W}_{a,out} &= \dot{m}(h_1 - h_{2a}) \end{split}$$

$$2 MW\left(\frac{1000 kJ}{1 MW}\right) = \dot{m}(3231.7 - 2682.4)\frac{kJ}{kg}$$

 $\dot{m} = 3.64 \, kg/s$ 

# **CLASS ACTIVITY**

- Air at 200 kPa and 950 K enters an adiabatic nozzle at low velocity and is discharged at a pressure of 110 kPa. If the isentropic efficiency of the nozzle is 92 percent, determine (assume constant specific heats of air):
  - a) The maximum possible exit velocity
  - b) The exit temperature
  - c) The actual exit velocity of the air

- Solution (assumptions):
  - □ Steady operating conditions exist
  - □ Air is an ideal gas
  - □ The kinetic and potential energies are negligible





• Solution (a):

$$\frac{T_{2s}}{T_1} = \left(\frac{P_{2s}}{P_1}\right)^{\frac{k-1}{k}}$$

$$T_{2s} = T_1 \left(\frac{P_{2s}}{P_1}\right)^{\frac{k-1}{k}} = (950 \text{ K}) \left(\frac{110 \text{ kPa}}{200 \text{ kPa}}\right)^{\frac{(0.349)}{1.349}} = 814 \text{ K}$$

• Solution (a):

 $e_{in} = e_{out}$ 

$$h_1 + \frac{V_1^2}{2} = h_{2s} + \frac{V_{2s}^2}{2}$$

$$V_{2s} = \sqrt{2(h_1 - h_{2s})} = \sqrt{2c_{p,avg}(T_1 - T_2)} = \sqrt{2(1.11\frac{kJ}{kg - K})(950 - 814)K(\frac{1000\frac{m^2}{s^2}}{1 kJ/kg})}$$

$$V_{2s} = 549 \frac{m}{s}$$

• Solution (b):

$$\eta_N = \frac{Actual \ KE \ nozzle \ exit}{Isentropic \ KE \ at \ nozzle \ exit} = \frac{V_{2a}^2}{V_{2s}^2}$$

$$0.92 = \frac{950 - T_{2a}}{950 - 814} \quad \rightarrow \quad T_{2a} = 825 \ K$$



• Solution (c):

$$\eta_N = \frac{Actual \ KE \ nozzle \ exit}{Isentropic \ KE \ at \ nozzle \ exit} = \frac{V_{2a}^2}{V_{2s}^2}$$

$$V_{2a} = \sqrt{\eta_N \times V_{2s}^2} = \sqrt{(0.92)(549\frac{m}{s})^2} = 527\frac{m}{s}$$

## **Chapter 8 Summary**

• We did not cover 8-12 (Entropy Balance)