## CAE 208 / MMAE 320: Thermodynamics Fall 2023

## November 21, 2023 <br> Entropy (3)

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Dr. Mohammad Heidarinejad, Ph.D., P.E. Civil, Architectural and Environmental Engineering Illinois Institute of Technology muh182@iit.edu

## ANNOUNCEMENTS

## Announcements

- Exam 2 solutions are posted
- Assignment 9 is posted (2-3 days more extension is fine to submit it on Friday or the weekend after the Thanksgiving)
- I will post assignment 10 due 12/01/23 (those who need to submit it)
- The bonus activities document is posted. Please pay attention to the deadlines and also the updates about this task (an idea submission by the 11/29/2023 is required)

RECAP

## Recap

- The equality in the Clausius inequality holds for totally or jut internally reversible cycles and the inequality for the irreversible ones:

$$
\left(\oint \frac{\delta Q}{T}\right)_{\text {int }, \text { rev }}=0
$$

## Recap

- A special case: Internally reversible isothermal heat transfer processes:

$$
\Delta S=\int_{1}^{2}\left(\frac{\delta Q}{T}\right)_{i n t, r e v}=\int_{1}^{2}\left(\frac{\delta Q}{T_{0}}\right)_{i n t, r e v}=\frac{1}{T_{0}} \int_{1}^{2} \delta Q_{i n t, r e v}
$$

$$
\Delta S_{\text {isothermal }}=\frac{Q}{T_{0}} \quad\left(\frac{k J}{K}\right)
$$



## Recap

- Entropy is a property:



## Recap

- For entropy, we can say "the increase of entropy principle":

$$
\begin{aligned}
& \oint_{1}^{2} \frac{\delta Q}{T} \leq \Delta S \\
& \Delta S_{s y s}=S_{2}-S_{1}=\oint_{1}^{2} \frac{\delta Q}{T}+S_{g e n}
\end{aligned}
$$

$\Delta S_{\text {isolated }} \geq 0$

## Recap

- For an isolated system:

$$
S_{\text {gen }}=\Delta S_{\text {total }}=\Delta S_{\text {sys }}+\Delta S_{\text {surr }} \geq 0
$$

$$
S_{g e n}=\left\{\begin{array}{cc}
>0 . & \text { irreversible proces } \\
=0 \quad \text { reversible process } \\
<0 . & \text { impossible process }
\end{array}\right.
$$



## Recap

- The entropy of a fixed mass can be changed by: - Heat Transfer
$\square$ Irreversibilities
- Entropy of a fixed mass does not change during a process that is internally reversible and adiabatic. During this process entropy remains constant and we call it isentropic process

$$
\Delta s=0 \text { or } s_{2}=s_{1} \quad\left(\frac{k J}{k g-K)}\right.
$$



## Recap

- We can rearrange our entropy equation:

$$
\delta Q_{i n t, r e v}=T d S
$$

$$
Q_{i n t, r e v}=\int_{1}^{2} T d S \quad(k J)
$$


(The area under the process curve on a T-S diagram represents heat transfer during an internally reversible process)

## WHAT IS ENTROPY

## What Is Entropy

- Entropy can be viewed as a measure of molecular disorder or molecular randomness (lowest in solid phase and highest in the gas phase)



## What Is Entropy

- Oscillations in solid phase fade as temperature is decreased, and the molecules supposedly become motionless at absolute zero. This represents a state of ultimate molecular order (and minimum energy).
- The entropy of a pure crystalline substance at absolute zero temperature is zero since there is no uncertainty about the state of the molecules at that instant (3rd law of thermodynamics)



## What Is Entropy

- The third law of thermodynamics provide an absolute reference point for the determination of entropy
- The entropy determined relative to this point is called absolute entropy and it is extremely useful in the thermodynamics analysis of entropy
- The entropy of a substance that is not pure crystalline (such as a solid solution) is not zero at absolute zero temperature


## What Is Entropy

- A few examples:



## THE T DS RELATIONS

## The T ds Relations

- The first T ds (or Gibbs) equation (closed stationary system):

$$
\begin{align*}
& \delta Q_{\text {int,rev }}-\delta W_{\text {int,rev,out }}=d U \\
&  \tag{kJ}\\
& \begin{array}{l}
T d S=d U+P d V \\
\\
\text { int,rev }
\end{array}=T d s \\
& \\
& \delta W_{\text {int }, \text { rev,out }}=P d V \\
& \\
& \\
& \\
& \left.\begin{array}{l}
T d s=d u+P d v \quad\left(\frac{k J}{k g}\right) \\
h=d u+P d v \\
h=u+P v \rightarrow d h=d u+P d v+v d P
\end{array}\right\} \rightarrow T d s=d h-v d P
\end{align*}
$$

## The T ds Relations

- Use the first T ds (or Gibbs) equation to solve for entropy changes

$$
\begin{aligned}
& d s=\frac{d u}{T}+\frac{P d v}{T} \\
& d s=\frac{d h}{T}-\frac{v d P}{T}
\end{aligned}
$$



$$
\begin{aligned}
& T d s=d u+P d \cup \\
& T d s=d h-\cup d P
\end{aligned}
$$

## ENTROPY CHANGE OF LIQUIDS AND SOLIDS

## Entropy Change of Liquids and Solids

- Liquids and solids can be approximated as incompressible substances ( $d v \cong 0 \& c_{p}=c_{v}=c_{p}=c$ ):

$$
\begin{aligned}
& d s=\frac{d h}{T}-\frac{v d P}{T} \\
& s_{2}-s_{1}=\int_{1}^{2} c(T) \frac{d T}{T} \cong c_{\text {avg }} \ln \left(\frac{T_{2}}{T_{1}}\right) \\
& s_{2}-s_{1}=\int_{1}^{2} c(T) \frac{d T}{T} \cong c_{\text {avg }} \ln \left(\frac{T_{2}}{T_{1}}\right)=0 \quad \rightarrow \quad T_{2}=T_{1}
\end{aligned}
$$

(For isentropic)

## THE ENTROPY CHANGE OF IDEAL GASES

## The Entropy Change of Ideal Gases

- For gases, we can write:

$$
\begin{aligned}
& d s=\frac{d u}{T}+\frac{P d v}{T}=c_{v} \frac{d T}{T}+R \frac{d v}{v} \\
& s_{2}-s_{1}=\int_{1}^{2} c_{v}(T) \frac{d T}{T}+R \ln \left(\frac{v_{2}}{v_{1}}\right) \\
& s_{2}-s_{1}=\int_{1}^{2} c_{p}(T) \frac{d T}{T}+R \ln \left(\frac{P_{2}}{P_{1}}\right)
\end{aligned}
$$

## The Entropy Change of Ideal Gases

- Approach 1: Constant Specific Heats (Approximate Analysis):

$$
\begin{gathered}
s_{2}-s_{1}=\int_{1}^{2} c_{v}(T) \frac{d T}{T}+R \ln \left(\frac{v_{2}}{v_{1}}\right) \\
s_{2}-s_{1}=c_{v, a v g} \ln \left(\frac{T_{2}}{T_{1}}\right)+R \times \ln \left(\frac{v_{2}}{v_{1}}\right)
\end{gathered}
$$



$$
\left\{\begin{array}{rl}
d h & =c_{p} d T \\
v & =\frac{R T}{P}
\end{array} \rightarrow s_{2}-s_{1}=c_{p, a v g} \times \ln \left(\frac{T_{2}}{T_{1}}\right)-R \times \ln \left(\frac{P_{2}}{P_{1}}\right)\right.
$$

## The Entropy Change of Ideal Gases

- Approach 2: Variable Specific Heats (Exact Analysis):

$$
\begin{aligned}
& s^{0}=\int_{0}^{T} c_{p}(T) \frac{d T}{T} \\
& \int_{0}^{T} c_{p}(T) \frac{d T}{T}=s_{2}^{0}-s_{1}^{0}
\end{aligned}
$$

$$
s_{2}-s_{1}=s_{2}^{0}-s_{1}^{0}-R \times \ln \left(\frac{P_{2}}{P_{1}}\right)
$$

$$
\overline{s_{2}}-\overline{s_{1}}=\overline{s_{2}^{0}}-\overline{s_{1}^{0}}-R_{u} \times \ln \left(\frac{P_{2}}{P_{1}}\right)
$$

## The Entropy Change of Ideal Gases

- Approach 2: Variable Specific Heats (Exact Analysis):



## CLASS ACTIVITY

## Class Activity

- Air is compressed from an initial state of 100 kPa and $17^{\circ} \mathrm{C}$ to a final state of 600 kPa and $57^{\circ} \mathrm{C}$. Determine the entropy change of air during this compression process by using:
a) Property values from the air table
b) Average specific heats


## Class Activity

- Solution (assumptions):
$\square$ Air is an ideal gas since it is at a high temperature and low pressure relative to its critical point values



## Class Activity

- Solution (calculations):
- Part (a): Table A-21
$s_{2}-s_{1}=s_{2}^{0}-s_{1}^{0}-R \times \ln \left(\frac{P_{2}}{P_{1}}\right)=\left[(1.79783-1.66802) \frac{k J}{k g-K}\right] \times \ln \left(\frac{600 k P a}{100 k P a}\right)$
$s_{2}-s_{1}=-0.3844 \frac{\mathrm{~kJ}}{\mathrm{~kg}-\mathrm{K}}$
- Part (b): Using a $c_{p}$ value at the average temperature of $37^{\circ} \mathrm{C}$ (Table A-2b)

$$
\begin{aligned}
& s_{2}-s_{1}=c_{p, a v g} \ln \left(\frac{T_{2}}{T_{1}}\right)-R \times \ln \left(\frac{P_{2}}{P_{1}}\right) \\
& =\left(1.006 \frac{\mathrm{~kJ}}{\mathrm{~kg}-\mathrm{K}}\right) \ln \left(\frac{330 \mathrm{~K}}{290 \mathrm{~K}}\right)-\left(0.287 \frac{\mathrm{~kJ}}{\mathrm{~kg}-\mathrm{K}}\right) \ln \left(\frac{600 \mathrm{kPa}}{100 \mathrm{kPa}}\right)=-0.3842 \frac{\mathrm{~kJ}}{\mathrm{~kg}-\mathrm{K}}
\end{aligned}
$$

## ISENTROPIC PROCESSES OF IDEAL GASES

## Isentropic Processes of Ideal Gases

- Approach 1: Constant Specific Heats (Approximate Analysis) for Isentropic Processes of Ideal Gases

$$
\begin{aligned}
& s_{2}-s_{1}=0=c_{v} \ln \left(\frac{T_{2}}{T_{1}}\right)+R \times \ln \left(\frac{v_{2}}{v_{1}}\right) \\
& \ln \left(\frac{T_{2}}{T_{1}}\right)=-\frac{R}{c_{v}} \ln \left(\frac{v_{2}}{v_{1}}\right) \\
& \ln \left(\frac{T_{2}}{T_{1}}\right)=\ln \left(\frac{v_{1}}{v_{2}}\right)^{\frac{R}{c_{v}}} \\
& \frac{T_{2}}{T_{1}}=\left(\frac{v_{1}}{v_{2}}\right)^{k-1}
\end{aligned}
$$

## Isentropic Processes of Ideal Gases

- Approach 1: Constant Specific Heats (Approximate Analysis) for Isentropic Processes of Ideal Gases

$$
\begin{aligned}
& \left(\frac{T_{2}}{T_{1}}\right)_{s=\text { constant }}=\left(\frac{v_{1}}{v_{2}}\right)^{k-1} \\
& \left(\frac{T_{2}}{T_{1}}\right)_{s=\text { constant }}=\left(\frac{P_{2}}{P_{1}}\right)^{\frac{k-1}{k}} \\
& \left(\frac{P_{2}}{P_{1}}\right)_{s=\text { constant }}=\left(\frac{v_{1}}{v_{2}}\right)^{k}
\end{aligned}
$$

## Isentropic Processes of Ideal Gases

- Approach 1: Constant Specific Heats (Approximate Analysis) for Isentropic Processes of Ideal Gases
$T v^{k-1}=$ Constant

$$
T P^{\frac{1-k}{k}}=\text { Constant }
$$

$$
P v^{k}=\text { Constant }
$$

$$
\begin{aligned}
&\left(\frac{T_{2}}{T_{1}}\right)_{s=\text { const. }}=\left(\frac{P_{2}}{P_{1}}\right)^{(k-1) / k}=\left(\frac{U_{1}}{U_{2}}\right)^{k-1} \\
& \text { *ideal gas } \\
& \text { Valid for } \text { *isentropic process } \\
& \text { *constant specific heats }
\end{aligned}
$$

## The Entropy Change of Ideal Gases

- Approach 2: Variable Specific Heats (Exact Analysis) for Isentropic Processes of Ideal Gases

$$
\begin{aligned}
& 0=s_{2}^{0}-s_{1}^{0}-R \times \ln \left(\frac{P_{2}}{P_{1}}\right) \\
& s_{2}^{0}=s_{1}^{0}+R \times \ln \left(\frac{P_{2}}{P_{1}}\right) \\
& s_{2}^{0}=s_{1}^{0}+R \times \ln \left(\frac{P_{2}}{P_{1}}\right) \rightarrow \frac{P_{2}}{P_{1}}=\exp \left(\frac{s_{2}^{0}-s_{1}^{0}}{R}\right)
\end{aligned}
$$

## CLASS ACTIVITY

## Class Activity

- Air enters an isentropic turbine at 150 psia and $900^{\circ} \mathrm{F}$ through a $0.5 \mathrm{ft}^{2}$ inlet section with a velocity of $500 \mathrm{ft} / \mathrm{s}$. It leaves at 15 psia with a velocity of $100 \mathrm{ft} / \mathrm{s}$. Calculate the air temperature at the turbine exit and the power produced, in hp , by this turbine.


## Class Activity

- Solution (assumptions):
- Steady flow
$\square$ The process is isentropic (both reversible and adiabatic)
$\square$ Ideal gas with a constant specific heat



## Class Activity

- Solution (Tables):
- Table A-2Eb: @ $600^{\circ} \mathrm{F} \rightarrow c_{p}=0.250 \frac{\mathrm{Btu}}{\mathrm{lbm-R}}$ and $k=1.3777$

| Ideal-gas specific heats of various common gases (b) At various temperatures |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Temp., ${ }^{\circ} \mathrm{F}$ | $c_{p} \mathrm{Btu} / \mathrm{lbm} \cdot \mathrm{R}$ | $c_{u} \mathrm{Btu} / \mathrm{lbm} \cdot \mathrm{R}$ | $k$ | $c_{p} \mathrm{Btu} / \mathrm{lbm} \cdot \mathrm{R}$ | $c_{v} \mathrm{Btu} / \mathrm{lbm} \cdot \mathrm{R}$ | $k$ |
|  | Air |  |  | Carbon dioxide, $\mathrm{CO}_{2}$ |  |  |
| 40 | 0.240 | 0.171 | 1.401 | 0.195 | 0.150 | 1.300 |
| 100 | 0.240 | 0.172 | 1.400 | 0.205 | 0.160 | 1.283 |
| 200 | 0.241 | 0.173 | 1.397 | 0.217 | 0.172 | 1.262 |
| 300 | 0.243 | 0.174 | 1.394 | 0.229 | 0.184 | 1.246 |
| 400 | 0.245 | 0.176 | 1.389 | 0.239 | 0.193 | 1.233 |
| 500 | 0.248 | 0.179 | 1.383 | 0.247 | 0.202 | 1.223 |
| 600 | 0.250 | 0.182 | 1.377 | 0.255 | 0.210 | 1.215 |
| 700 | 0.254 | 0.185 | 1.371 | 0.262 | 0.217 | 1.208 |
| 800 | 0.257 | 0.188 | 1.365 | 0.269 | 0.224 | 1.202 |
| 900 | 0.259 | 0.191 | 1.358 | 0.275 | 0.230 | 1.197 |
| 1000 | 0.263 | 0.195 | 1.353 | 0.280 | 0.235 | 1.192 |
| 1500 | 0.276 | 0.208 | 1.330 | 0.298 | 0.253 | 1.178 |
| 2000 | 0.286 | 0.217 | 1.312 | 0.312 | 0.267 | 1.169 |

## Class Activity

- Solution (Tables):
- Table A-1E: $R=0.3704 \frac{p \text { psia- }-t^{3}}{l b m-R}$

| TABLE A-1E |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Molar mass, gas constant, and critical-point properties |  |  |  |  |  |  |  |
| Substance | Formula | Molar mass, $M$ lbm/lbmol | Gas constant, $R^{*}$ |  | Critical-point properties |  |  |
|  |  |  | Btu/lbm . | $\mathrm{psia} \cdot \mathrm{ft}^{3} / \mathrm{lbm} \cdot \mathrm{R}$ | Temperature, R | Pressure, psia | Volume, $\mathrm{ft}^{3} / \mathrm{lbmol}$ |
| Air | - | 28.97 | 0.06855 | 0.3704 | 238.5 | 547 | 1.41 |

## Class Activity

- Solution (Problem solving):

$$
\begin{aligned}
& \dot{m}=\dot{m}_{1}=\dot{m}_{2} \\
& \dot{E}_{\text {in }}-\dot{E}_{\text {out }}=\frac{d \dot{E}_{\text {system }}}{d t}=0
\end{aligned}
$$



$$
\begin{aligned}
& \dot{m}\left(h_{1}+V_{1}^{2}\right)=\dot{m}\left(h_{2}+\frac{V_{2}^{2}}{2}\right)+\dot{W}_{\text {out }} \\
& \dot{W}_{\text {out }}=\dot{m}\left(h_{1}-h_{1}+\frac{V_{1}^{2}-V_{2}}{2}\right)
\end{aligned}
$$

## Class Activity

- Solution (Calculations):

$$
\begin{aligned}
& \frac{T_{2}}{T_{1}}=\left(\frac{P_{2}}{P_{1}}\right)^{\frac{k-1}{k}} \rightarrow T_{2}=T_{1} \times\left(\frac{P_{2}}{P_{1}}\right)^{\frac{k-1}{k}}=(900+460 R)\left(\frac{15 p \text { sia }}{150 \text { psia }}\right)^{\frac{0.3777}{1.377}}=724 R \\
& v_{1}=\frac{R T_{1}}{P_{1}}=\frac{\left(0.3704 \frac{p \text { sia }-f t^{3}}{l b m-R}\right)(900+460 R)}{150 p s i a}=3.358 \frac{f t^{3}}{l b m} \\
& \dot{m}=\frac{A_{1} V_{1}}{v_{1}}=\frac{\left(0.5 f t^{2}\right)\left(500 \frac{f t}{s}\right)}{3.358 \frac{f t^{3}}{l b m}}=74.45 \frac{\mathrm{lbm}}{\mathrm{~s}}
\end{aligned}
$$

## Class Activity

- Solution (Calculations):

$$
\begin{aligned}
& \dot{W}_{\text {out }}=\dot{m}\left(h_{1}-h_{1}+\frac{V_{1}^{2}-V_{2}}{2}\right) \\
& \dot{W}_{\text {out }}=\left(74.45 \frac{l b m}{s}\right)\left[\left(0.250 \frac{B t u}{l b m-R}\right)(1360-724 R)+\left(\frac{\left(500 \frac{f t}{s}\right)^{2}}{2}-\frac{\left(100 \frac{f t}{s}\right)^{2}}{2}\right)\left(\frac{1 \frac{B t u}{l b m}}{\frac{25.037}{s^{2}} f t^{2}}\right)\right. \\
& \dot{W}_{\text {out }}=12,194 \frac{B t u}{s}\left(\frac{1 \mathrm{hp}}{0.7068 \frac{B t u}{s}}\right)=17,250 \mathrm{hp}
\end{aligned}
$$

## REVERSIBLE STEADY-FLOW WORK

## Reversible Steady-Flow Work

- Recall we had

$$
W_{b}=\int_{1}^{2} P d V
$$

- For steady flow:

$$
\begin{aligned}
& \delta q_{r e v}-\delta w_{r e v}=d h+d k e+d p e \\
& \left.\begin{array}{l}
\delta q_{r e v}=T d \\
T d s=d h-v d P
\end{array}\right\} \rightarrow \delta q_{r e v}=d h-v d P \\
& \qquad w_{r e v}=-\int_{1}^{2} v d P-\Delta k e-\Delta p e
\end{aligned}
$$

## Reversible Steady-Flow Work

- For steady flow:

$$
w_{r e v}=-\int_{1}^{2} v d P-\Delta k e-\Delta p e
$$


(a) Steady-flow system

(b) Closed system

## Reversible Steady-Flow Work

- We write the Bernoulli equation:

$$
w_{r e v}=0=v\left(P_{2}-P_{1}\right)+\left(\frac{V_{2}^{2}-V_{1}^{2}}{2}\right)+g\left(z_{2}-z_{1}\right)
$$

## ISENTROPIC EFFICIENCIES OF STEADY-FLOW DEVICES

## Isentropic Efficiencies

- Steady-flow devices deliver the most and consume the least work when the process is reversible:



## Isentropic Efficiencies

- Isentropic efficiency of a turbine can be written as:

$$
\eta_{T}=\frac{\text { Actual turbine work }}{\text { Isentropic turbine work }}=\frac{w_{a}}{w_{s}}
$$

$$
\eta_{T} \cong \frac{h_{1}-h_{2 a}}{h_{1}-h_{2 s}}
$$



## Isentropic Efficiencies

- Isentropic efficiency of compressors and pumps
$\eta_{C}=\frac{\text { Isentropic compressor work }}{\text { Actual compressor work }}=\frac{w_{s}}{w_{a}}$

$$
\eta_{C} \cong \frac{h_{2 s}-h_{1}}{h_{2 a}-h_{1}}
$$



## Isentropic Efficiencies

- Isentropic efficiency of nozzles

$$
\eta_{N}=\frac{\text { Actual } K E \text { nozzle exit }}{\text { Isentropic } K E \text { at nozzle exit }}=\frac{V_{2 a}^{2}}{V_{2 s}^{2}}
$$

$$
h_{1}=h_{2 a}+\frac{V_{2 a}^{2}}{2}
$$

$$
\eta_{N} \cong \frac{h_{1}-h_{2 a}}{h_{1}-h_{2 s}}
$$



## CLASS ACTIVITY

## Class Activity

- Steam enters an adiabatic turbine steadily at 3 MPa and 400 C and leaves at 50 kPa and 100 C . If the power output of the turbine is 2 MW , determine:
a) The isentropic efficiency of the turbine
b) The mass flow rate of the steam flowing through the turbine


## Class Activity

- Solution (assumptions):
$\square$ Steady operating conditions exist
The kinetic and potential energies are negligible




## Class Activity

- Solution (Tables):

State 1: $\quad\left\{\begin{array}{l}P_{1}=3 \mathrm{MPa} \\ T_{1}=400^{\circ} \mathrm{C}\end{array} \rightarrow\left\{\begin{array}{c}h_{1}=3231.7 \mathrm{~kJ} / \mathrm{kg} \\ s_{1}=6.9235 \mathrm{~kJ} /(\mathrm{kg}-\mathrm{K})\end{array}\right.\right.$

State 2a: $\quad\left\{\begin{array}{l}P_{1}=50 \mathrm{kPa} \\ T_{2 a}=100^{\circ} \mathrm{C}\end{array} \rightarrow \quad h_{2 a}=2682.4 \mathrm{~kJ} / \mathrm{kg}\right.$

State $2 s: \quad\left\{\begin{array}{c}P_{1}=50 \mathrm{kPa} \\ s_{1}=s_{2}\end{array} \rightarrow\left\{\begin{array}{l}s_{f}=1.0912 \mathrm{~kJ} /(\mathrm{kg}-\mathrm{K}) \\ s_{g}=7.5931 \mathrm{~kJ} /(\mathrm{kg}-\mathrm{K})\end{array}\right.\right.$

$$
\begin{aligned}
x_{2 s} & =\frac{s_{2 s}-s_{f}}{s_{f g}}=\frac{6.9235-1.0912}{6.5019}=0.897 \\
h_{2 s} & =h_{f}+x_{2 s} \times h_{f g}=340.54+0.897 \times(2304.7)=2407.9 \mathrm{~kJ} / \mathrm{kg}
\end{aligned}
$$

## Class Activity

- Solution (a): The isentropic efficiency is:

$$
\eta_{T}=\frac{h_{1}-h_{2 a}}{h_{1}-h_{2 s}}=\frac{3231.7-2682.4}{3231.7-2407.9}=0.667(\text { or } 66.7 \%)
$$

## Class Activity

- Solution (b): The mass flow rate is:

$$
\begin{aligned}
& \dot{E}_{\text {in }}=\dot{E}_{\text {out }} \\
& \dot{m} h_{1}=\dot{W}_{a, \text { out }}+\dot{m} h_{2 a} \\
& \dot{W}_{a, \text { out }}=\dot{m}\left(h_{1}-h_{2 a}\right) \\
& 2 M W\left(\frac{1000 \mathrm{~kJ}}{1 \mathrm{MW}}\right)=\dot{m}(3231.7-2682.4) \frac{\mathrm{kJ}}{\mathrm{~kg}} \\
& \dot{m}=3.64 \mathrm{~kg} / \mathrm{s}
\end{aligned}
$$

## CLASS ACTIVITY

## Class Activity

- Air at 200 kPa and 950 K enters an adiabatic nozzle at low velocity and is discharged at a pressure of 110 kPa . If the isentropic efficiency of the nozzle is 92 percent, determine (assume constant specific heats of air):
a) The maximum possible exit velocity
b) The exit temperature
c) The actual exit velocity of the air


## Class Activity

- Solution (assumptions):
- Steady operating conditions exist

Air is an ideal gas
$\square$ The kinetic and potential energies are negligible



## Class Activity

- Solution (a):

$$
\begin{aligned}
& \frac{T_{2 s}}{T_{1}}=\left(\frac{P_{2 s}}{P_{1}}\right)^{\frac{k-1}{k}} \\
& T_{2 s}=T_{1}\left(\frac{P_{2 s}}{P_{1}}\right)^{\frac{k-1}{k}}=(950 \mathrm{~K})\left(\frac{110 \mathrm{kPa}}{200 \mathrm{kPa}}\right)^{\frac{(0.349)}{1.349}}=814 \mathrm{~K}
\end{aligned}
$$

## Class Activity

- Solution (a):

$$
e_{\text {in }}=e_{\text {out }}
$$

$h_{1}+\frac{V_{1}^{2}}{2}=h_{2 s}+\frac{V_{2 s}^{2}}{2}$
$V_{2 s}=\sqrt{2\left(h_{1}-h_{2 s}\right)}=\sqrt{2 c_{p, a v g}\left(T_{1}-T_{2}\right)}=\sqrt{2\left(1.11 \frac{\mathrm{~kJ}}{\mathrm{~kg}-K}\right)(950-814) K\left(\frac{1000 \frac{\mathrm{~m}^{2}}{\mathrm{~s}^{2}}}{1 \mathrm{~kJ} / \mathrm{kg}}\right.}$
$V_{2 s}=549 \frac{\mathrm{~m}}{\mathrm{~s}}$

## Class Activity

- Solution (b):

$$
\begin{aligned}
& \eta_{N}=\frac{\text { Actual KE nozzle exit }}{\text { Isentropic KE at nozzle exit }}=\frac{V_{2 a}^{2}}{V_{2 s}^{2}} \\
& 0.92=\frac{950-T_{2 a}}{950-814} \quad \rightarrow \quad T_{2 a}=825 \mathrm{~K}
\end{aligned}
$$



## Class Activity

- Solution (c):

$$
\begin{aligned}
& \eta_{N}=\frac{\text { Actual KE nozzle exit }}{\text { Isentropic KE at nozzle exit }}=\frac{V_{2 a}^{2}}{V_{2 s}^{2}} \\
& V_{2 a}=\sqrt{\eta_{N} \times V_{2 s}^{2}}=\sqrt{(0.92)\left(549 \frac{\mathrm{~m}}{\mathrm{~s}}\right)^{2}}=527 \frac{\mathrm{~m}}{\mathrm{~s}}
\end{aligned}
$$

## Chapter 8 Summary

- We did not cover 8-12 (Entropy Balance)

